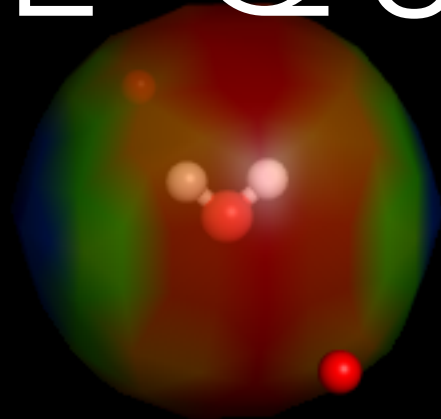


DAVID M. ROGERS, UNIV. SOUTH FLORIDA

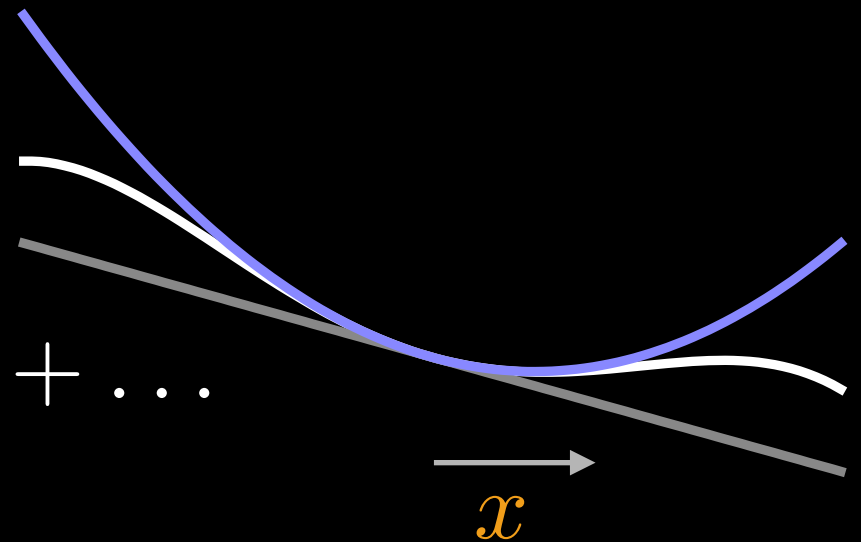
MULTIPOLES FROM REAL SPACE QUADRATURE

(SERM)ACS

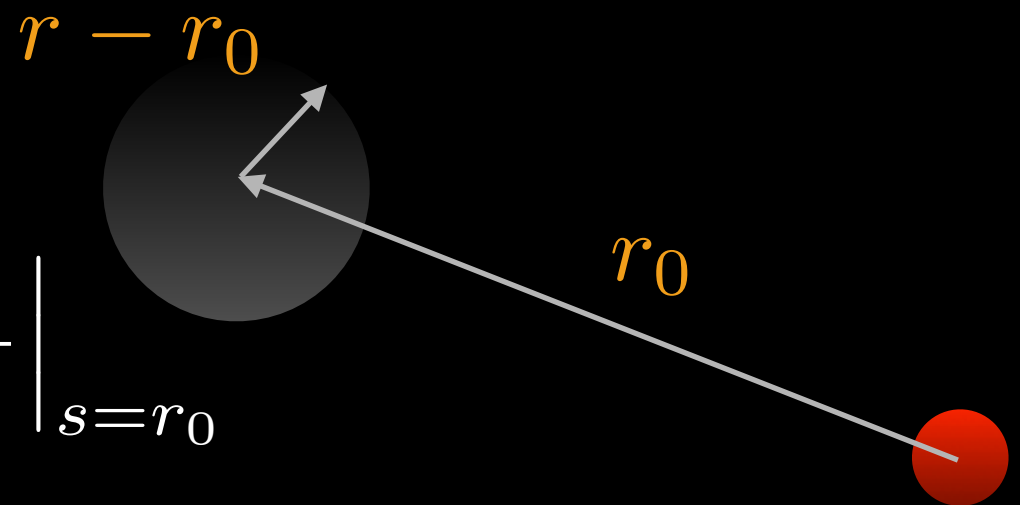
Nashville, 2014



$$f(x) = f(0) + x f'(0) + \frac{x^2}{2} f''(0) + \dots$$



$$\Phi(r) = \sum_{n=0} \frac{((r - r_0) \cdot \partial_s)^n}{n!} \frac{1}{|s|} \Big|_{s=r_0}$$



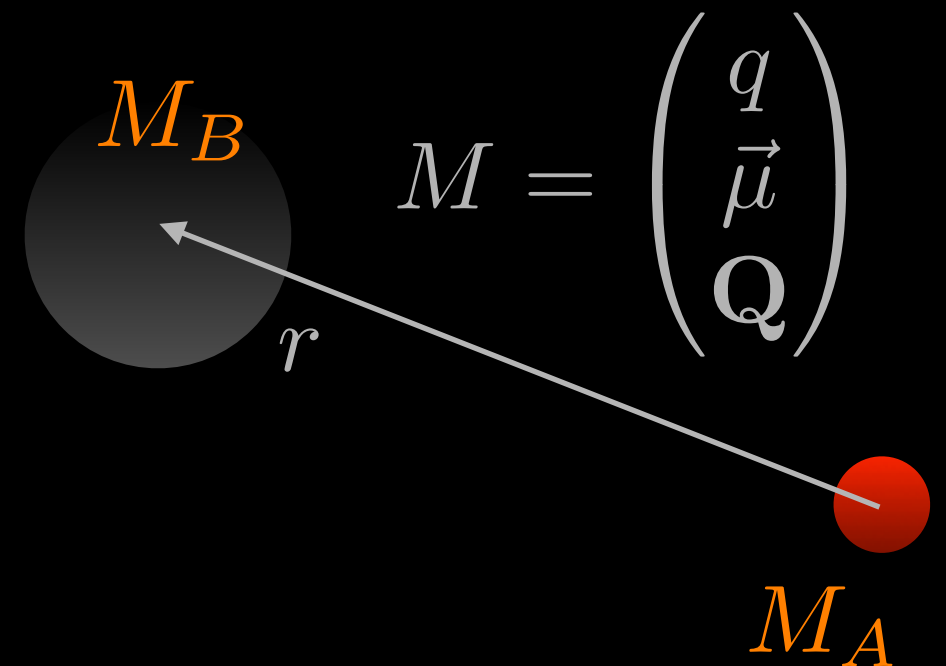
WHAT IS A MULTIPOLE EXPANSION?

A coefficient.

$$V_{AB}^{\text{es}} = M_A \cdot T_{AB} \cdot \tilde{M}_B$$

FAR-FIELD INTERACTION E

- Potential function expansion
 - Transferable, Polarizable FFs
 - Fast multipole method, $O(N)$



$$T = \begin{bmatrix} r^{-1} & \partial r^{-1} & \partial^{(2)} r^{-1} \\ \partial r^{-1} & \partial^{(2)} r^{-1} & \partial^{(3)} r^{-1} \\ \partial^{(2)} r^{-1} & \partial^{(3)} r^{-1} & \partial^{(4)} r^{-1} \end{bmatrix}$$

Jacobson, Williams, Herbert, JCP 130, 2009.

$\partial_{\alpha\beta\gamma\delta}^{(4)}$ Cartesian: $3^4 = 81$ component tensor!

$$M_{lm} = \sum_i q_i |r_p|^l C_{lm}(\hat{r}_p)$$

$$V_{l,k} = \binom{2l+2k}{2l}^{1/2} R^{-l-k-1} \sum_{m=-l-k}^{l+k} (-1)^m C_{l+k,-m}(\hat{R}) [M_A^{(l)} \otimes \tilde{M}_B^{(k)}]_{l+k,m}$$

$$[M_A^{(l)} \otimes \tilde{M}_B^{(k)}]_{l+k,m} = \sum_{n=-l}^l \sum_{j=-k}^k M_{A,n}^{(l)} \tilde{M}_{B,j}^{(k)} \langle l, n; k, j | l+k, m \rangle$$

Jeziorski, Moszynski, Szalewicz, Chem. Rev. 94, 1994.


BUT IT'S NOT SO SIMPLE?

Spherical multipoles are efficient, yet unintelligible.

THE REAL PROBLEM

$$\begin{aligned}\frac{1}{|x-y|} &= e^{-y \cdot \partial} \frac{1}{|x|} \\ &= \sum_{l=0}^{\infty} \frac{|y|^l}{|x|^{l+1}} P_l(\hat{x} \cdot \hat{y}) \\ &= \sum_{l=0}^{\infty} \sum_{m=-l}^l O_{l,m}(x) M_{l,m}(y)\end{aligned}$$

One step too far!



$$O_{lm}(x) = |x|^l (l + |m|)! P_{lm}(\cos \theta_x) e^{-im\phi_x}$$

$$M_{lm}(x) = |x|^{-l-1} (l - |m|)!^{-1} P_{lm}(\cos \theta_x) e^{im\phi_x}$$

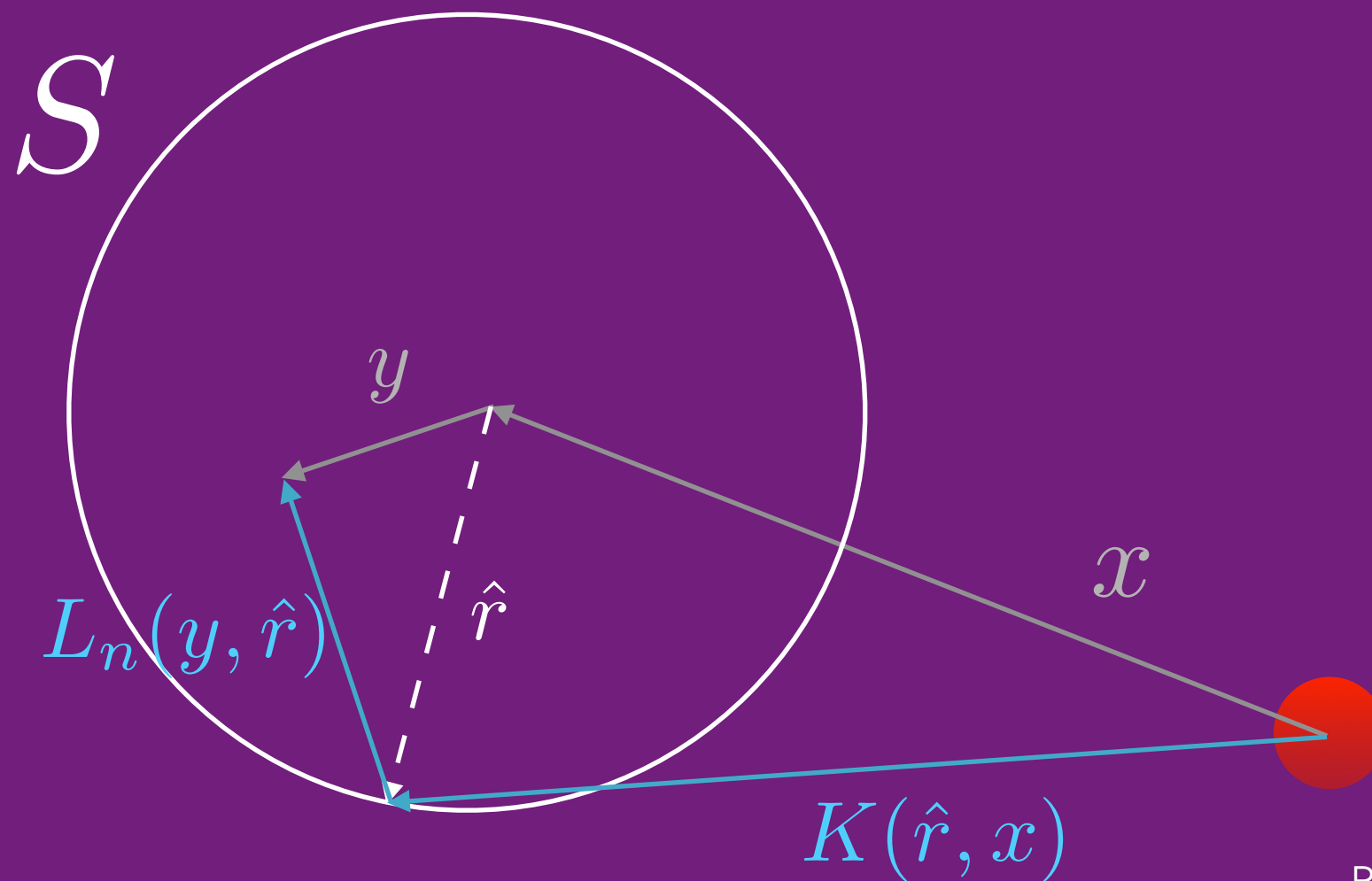
THE REAL PROBLEM

$$\begin{aligned}
 \frac{1}{|x - y|} &= e^{-y \cdot \partial} \frac{1}{|x|} \\
 &= \sum_{l=0} \frac{|y|^l}{|x|^{l+1}} P_l(\hat{x} \cdot \hat{y}) \\
 &= \sum_{l=0} \frac{|y|^l}{|x|^{l+1} l!} \hat{x}^{(l)} \hat{y}^{(l)}
 \end{aligned}$$

One step too far!



$$\partial^{(n)} |r|^{-1} = (-1)^n |r|^{-2n-1} \mathcal{T}_n r^{(n)}$$



Rogers (submitted), 2014.

$$(-y \cdot \partial)^n |x|^{-1} / n! = L_n(y, x) \quad \text{directional derivative}$$

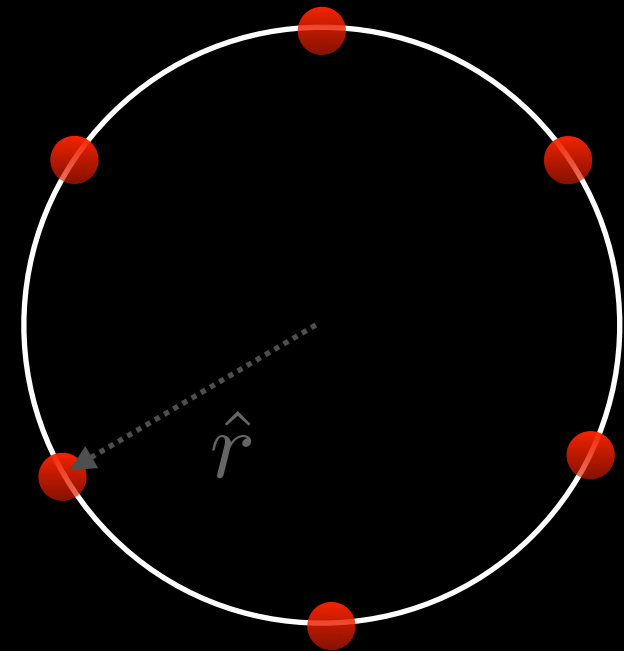
$$= \int_S L_n(y, \hat{r}) K(\hat{r}, x) d^2 \hat{r} \quad \text{spherical projection}$$

BACK TO THE DRAWING BOARD.

Ahrens and Beylkin, Proc. Royal Soc. A, 465, 2009.

QUADRATURE

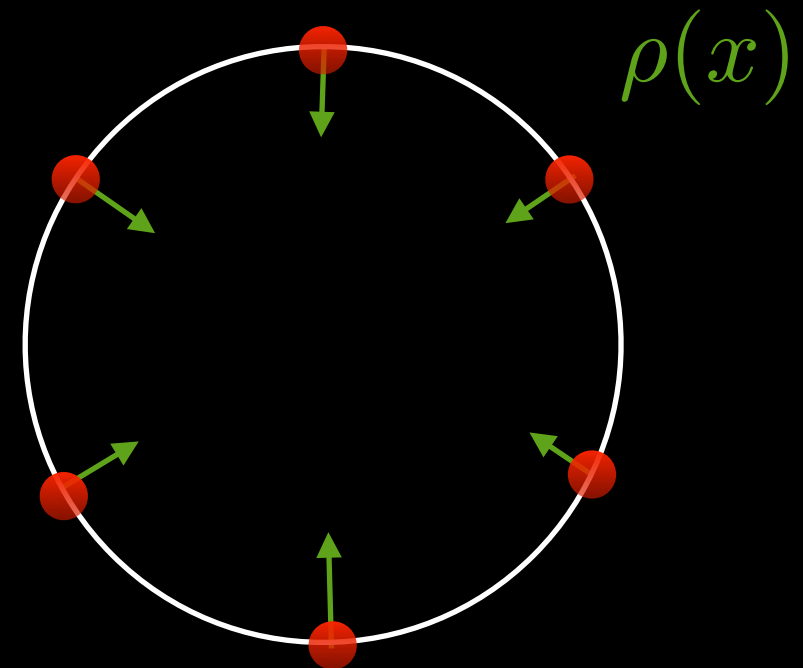
- Optimal number of points (p^2)
- Physical interpretation = surface charge distribution



$$\int_S f(y, \hat{r}) \sigma(\hat{r}) d^2 \hat{r}$$
$$= \sum_i \underbrace{(w_i \sigma(\hat{r}_i))}_{q_i} f(y, \hat{r}_i)$$

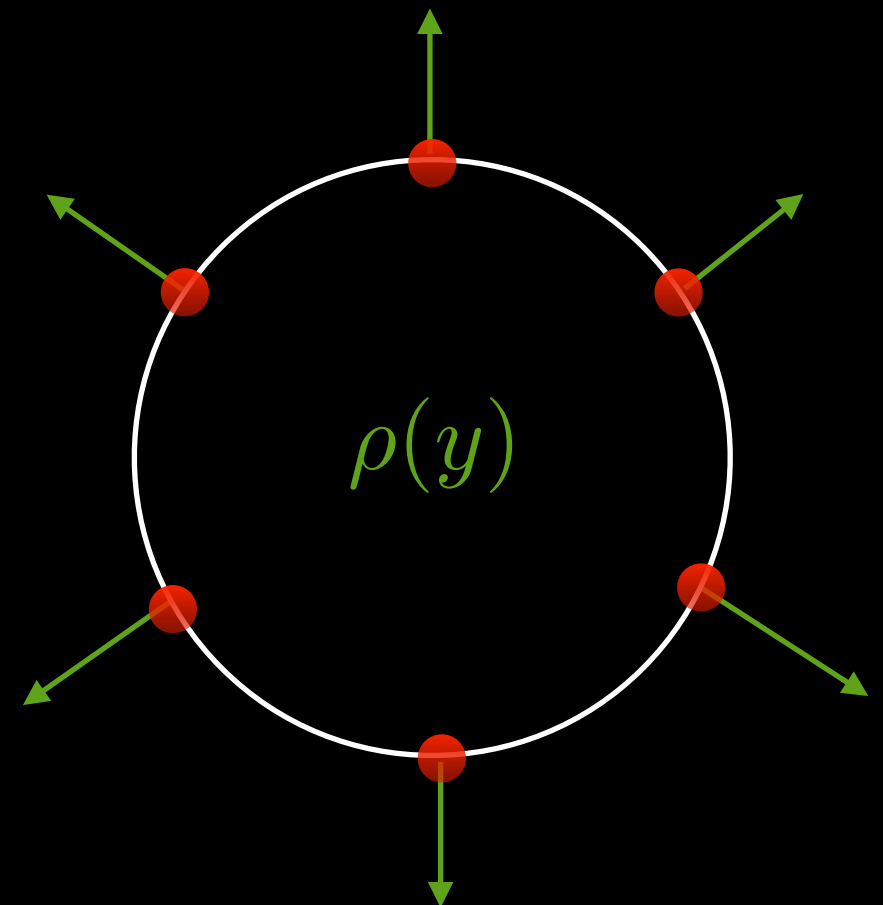
INNER EXPANSION

$$\sigma_i(\hat{r}) = \int K(\hat{r}, x) \rho(x) d^3x$$



OUTER EXPANSION

$$\sigma_o(\hat{r}) = \int \rho(y) K(y, \hat{r}) d^3y$$

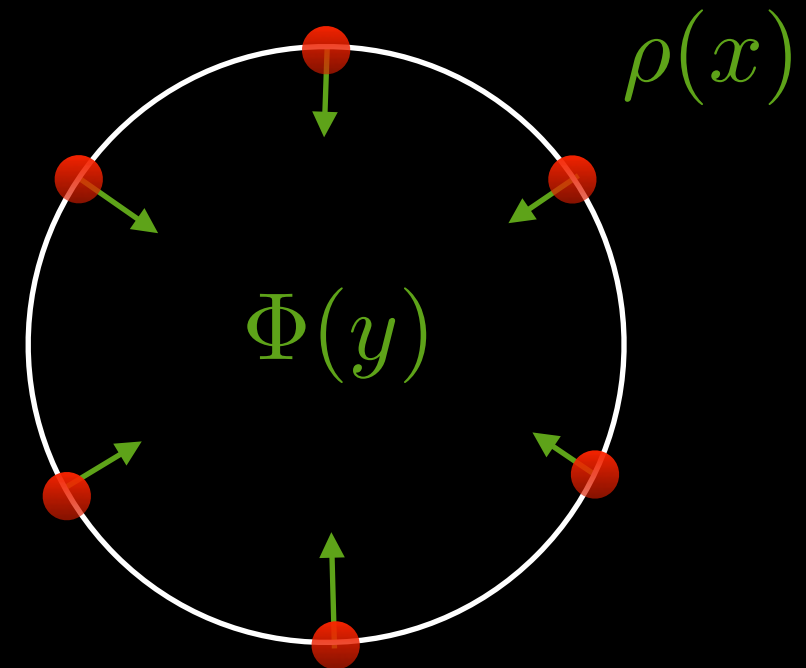


- Getting in is easy

INNER EXPANSION

$$\Phi(y) = \int_S \sum_n L_n(y, \hat{r}) \sigma_i(\hat{r}) d^2 \hat{r}$$

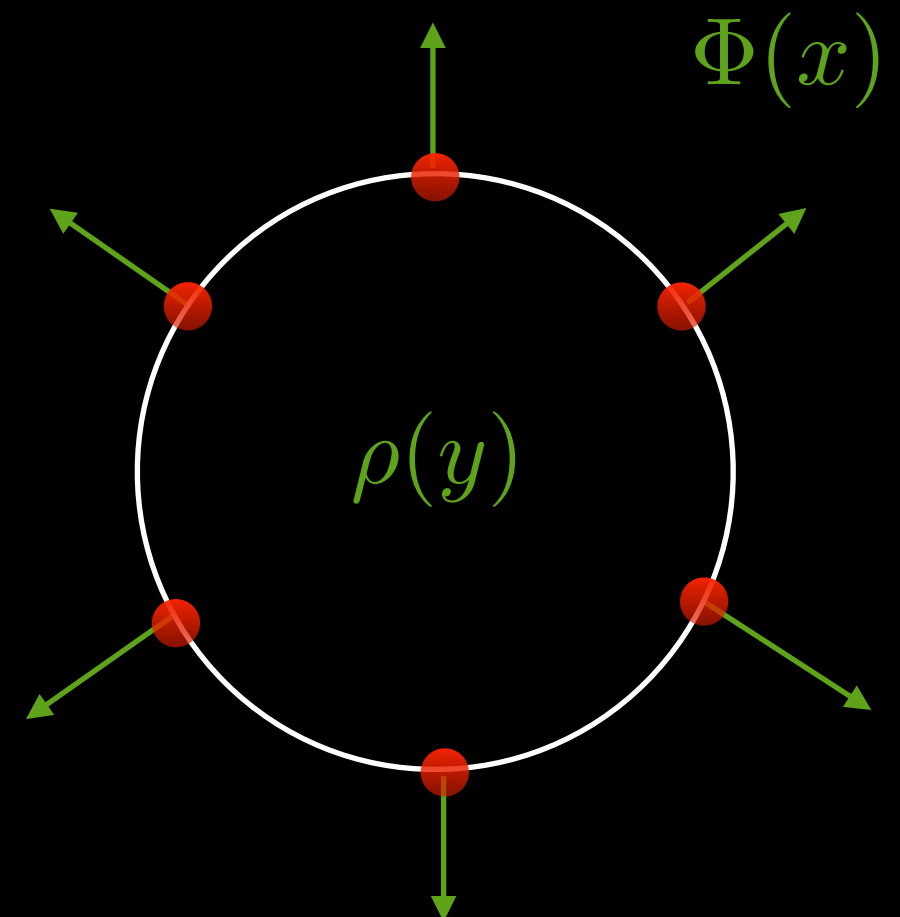
$$\simeq \frac{q_i}{|y - \hat{r}_i|} \quad (\text{may require scaling } S)$$



OUTER EXPANSION

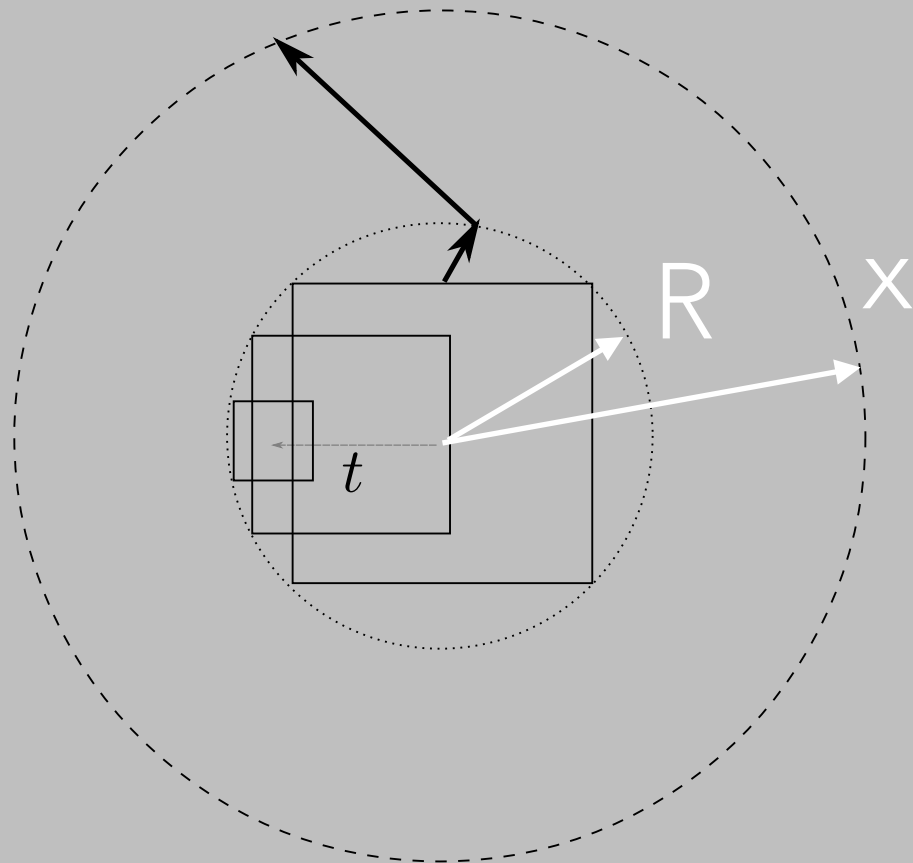
$$\Phi(x) = \int_S \sigma_o(\hat{r}) \sum_n L_n(\hat{r}, x) d^2 \hat{r}$$

$$\simeq \frac{q_i}{|x - \hat{r}_i|}$$

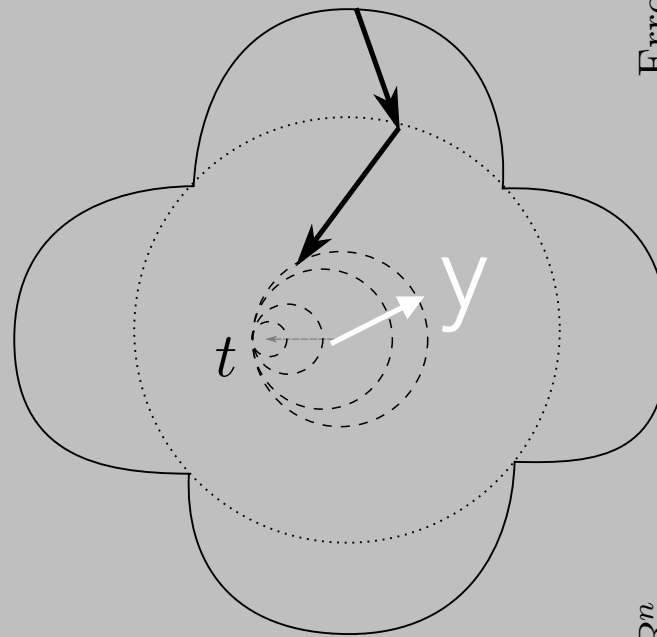


- Getting out is **even easier!**

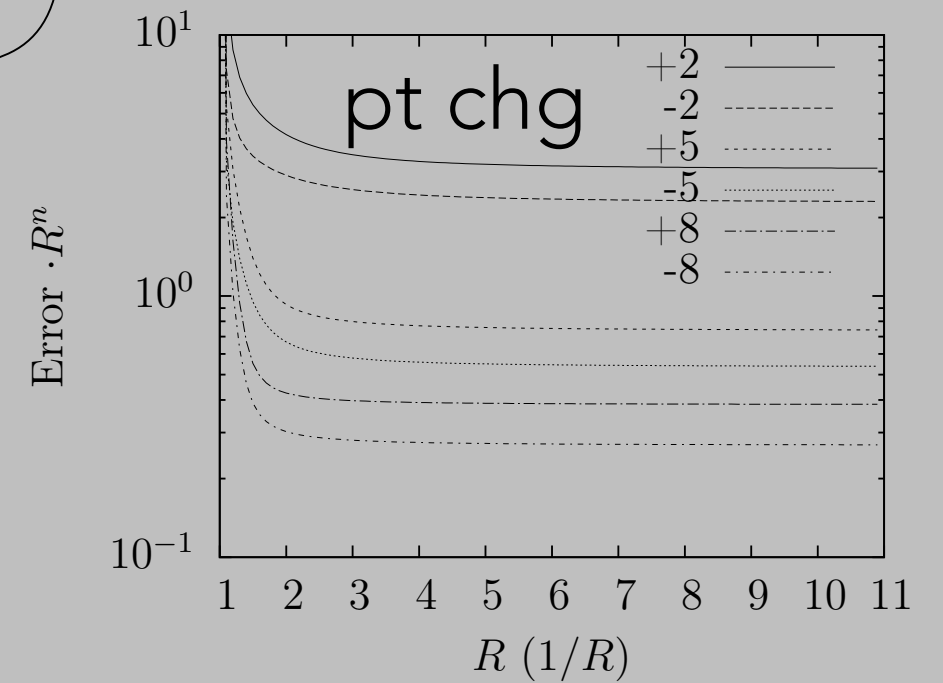
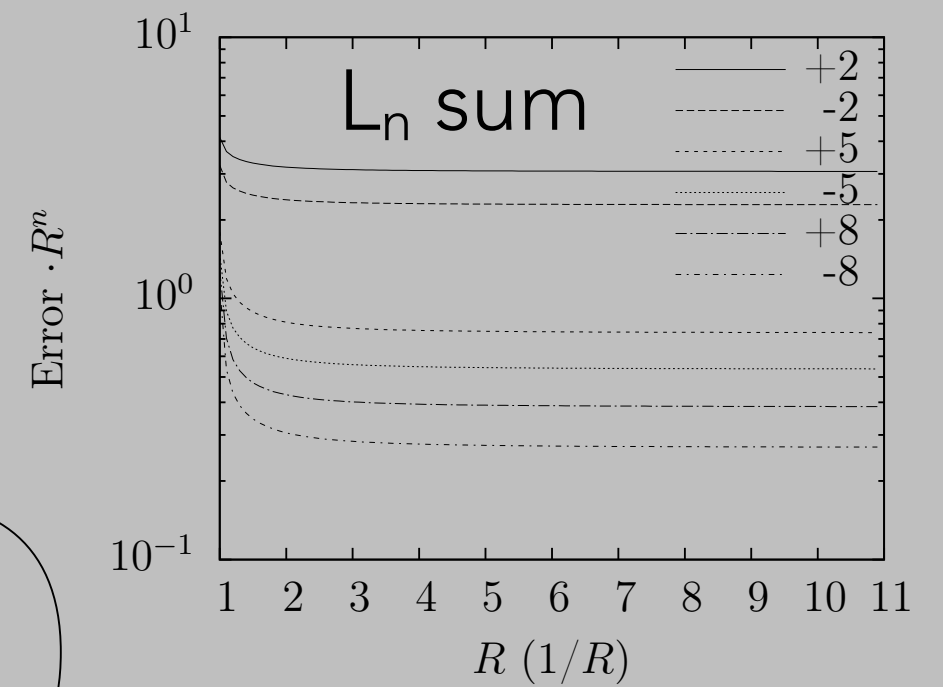
ERROR?



$$O\left(\frac{R}{x}\right)^{p+1}$$

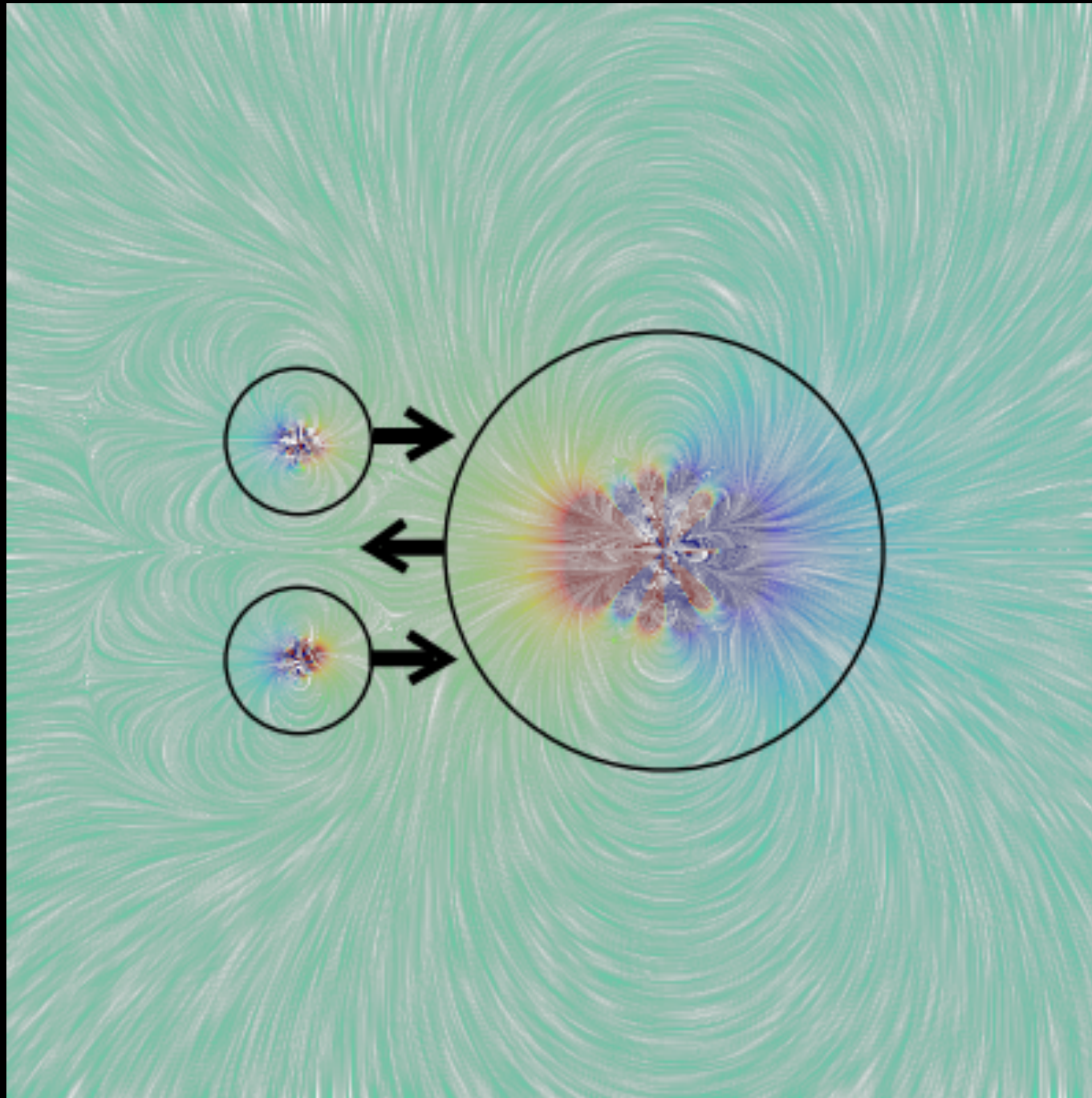


$$O\left(\frac{y}{R}\right)^p$$



Prefactor (au)

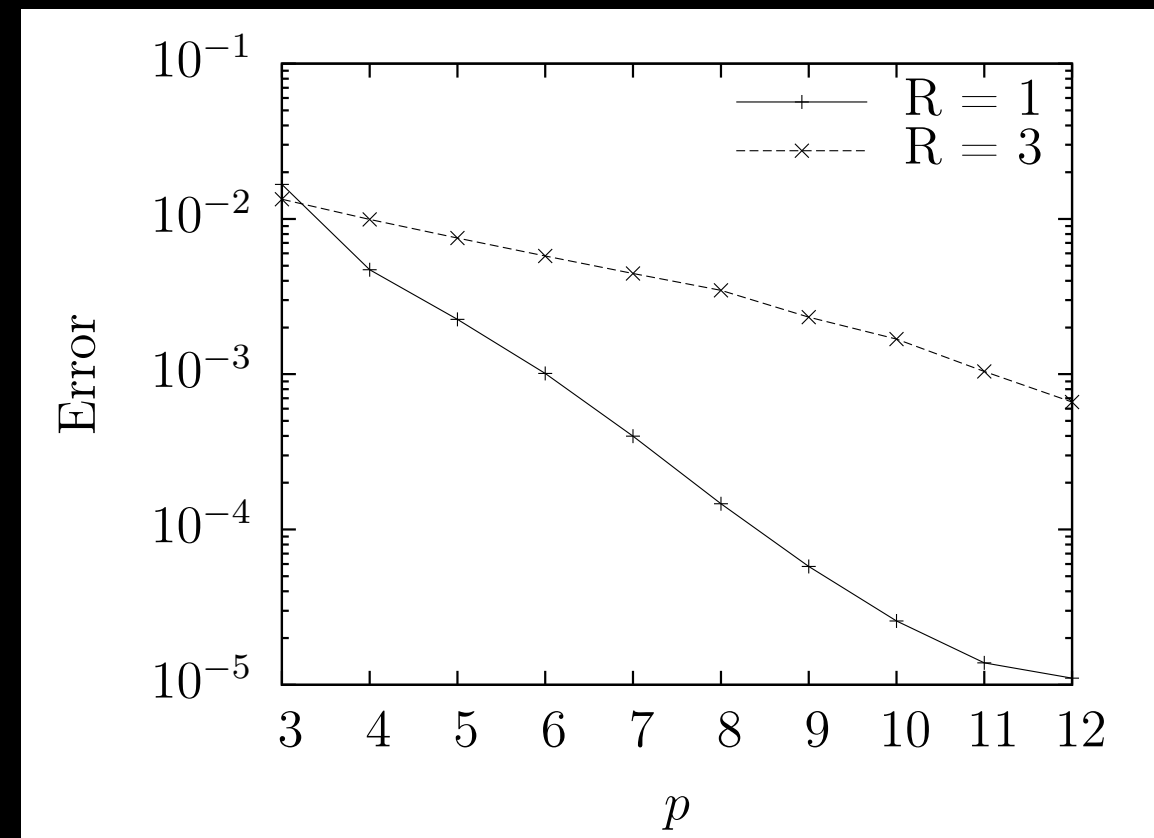
Rogers (submitted), 2014.



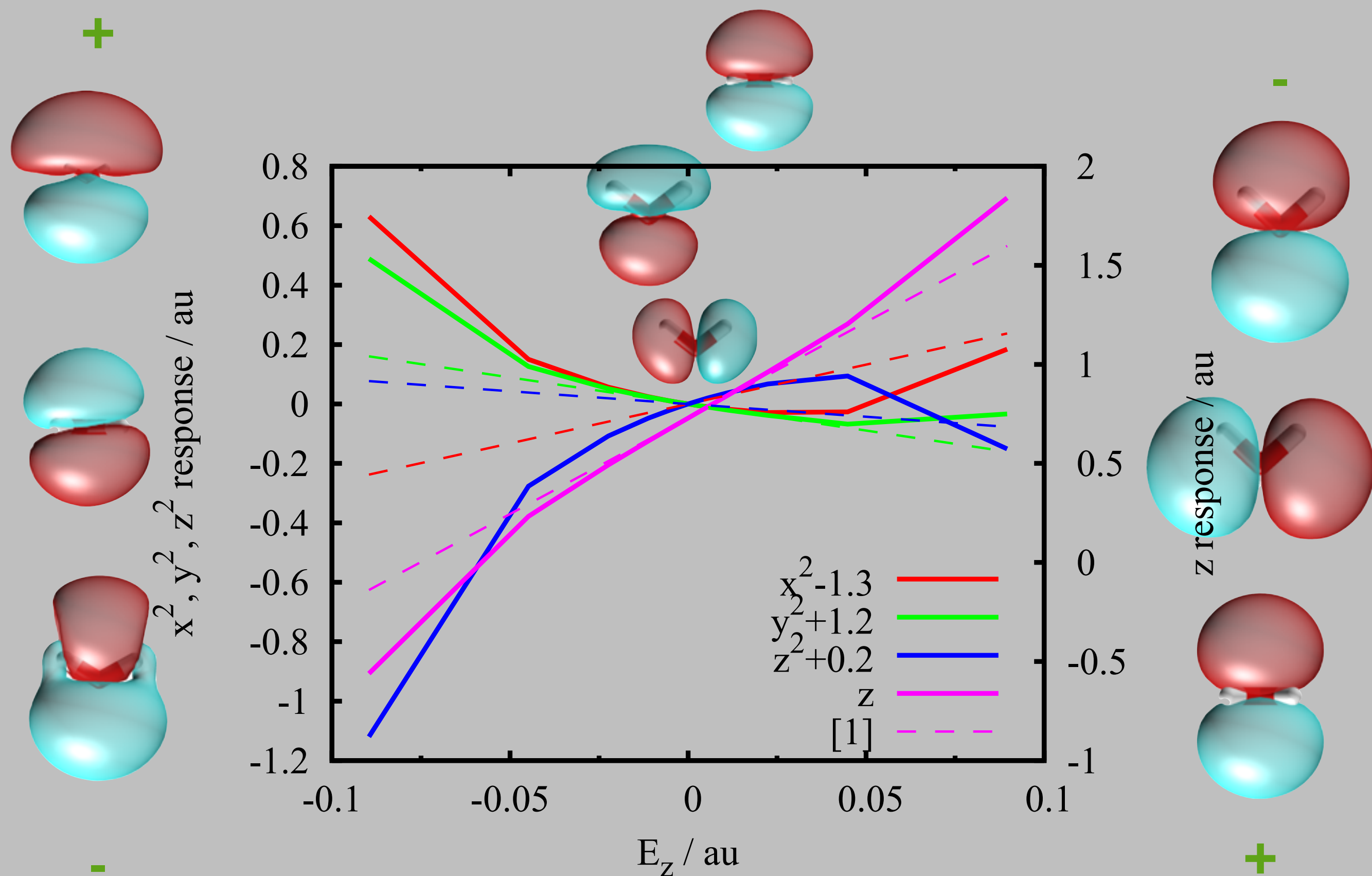
$$v(x) \equiv -\partial\Phi(x)$$

$$0 = \partial^2\Phi(x)$$

$$n \cdot v(x) = v_0(x), x \in \partial\Omega$$



FUN WITH MULTIPOLES



WORK WITH MULTIPOLES

[1] Elking et. al., J. Comput. Chem. 32, 2011.

CONCLUSIONS

- Spherical harmonics considered harmful
 - May have convergence issues!
Makino, J. Comp. Phys. 151, 2009.
- Real-space Quadrature
 - Basis vectors are real
 - Weights are charges
 - Approx. power and dimension identical
 - Symmetry is apparent
 - Shifting functions are identical to initial fitting!
Rogers (submitted), 2014.

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